

# On integrability of one third-order nonlinear evolution equation

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## Abstract

We study one third-order nonlinear evolution equation, recently introduced by Chou and Qu in a problem of plane curve motions, and find its transformation to the modified Korteweg–de Vries equation, its zero-curvature representation with an essential parameter, and its second-order recursion operator.

## 1 Introduction

In their recent work on the motions of plane curves [1], Chou and Qu found the following new third-order nonlinear evolution equation:

$$u_t = \frac{1}{2} \left( (u_{xx} + u)^{-2} \right)_x. \quad (1)$$

“We do not know if this equation arises from the AKNS- or the WKI-scheme”, wrote Chou and Qu in [1].

In the present paper, we study integrability of (1). In Section 2, we find a chain of Miura-type transformations, which relates the equation (1) with the modified Korteweg–de Vries equation (mKdV). In Section 3, using the obtained transformations and the well-known zero-curvature representation (ZCR) of the mKdV, we derive a complicated nontrivial ZCR of (1), which turns out to be neither AKNS- nor WKI-type one; and then we prove that simpler ZCRs of the equation (1) are trivial. In Section 4, we derive a second-order recursion operator of (1) from the obtained ZCR. Section 5 gives some concluding remarks.

## 2 Transformation to the mKdV

Let us try to transform the equation (1) into one of the well-known integrable equations. We do it, following the way described in [2]; further details on Miura-type transformations of scalar evolution equations can be found in [3].

First, we try to relate the equation (1) with an evolution equation of the form

$$v_t = v^3 v_{xxx} + g(v, v_x, v_{xx}) \quad (2)$$

by a transformation of the type

$$v(x, t) = a(u, u_x, \dots, u_{x\dots x}). \quad (3)$$

If there exists a transformation (3) between an evolution equation

$$u_t = f(u, u_x, u_{xx}, u_{xxx}) : \quad \partial f / \partial u_{xxx} \neq \text{constant} \quad (4)$$

and an equation of the form (2), then necessarily

$$a = (\partial f / \partial u_{xxx})^{1/3}. \quad (5)$$

Applying the transformation

$$(x, t, u(x, t)) \mapsto (x, t, v(x, t)) : \quad v = -(u_{xx} + u)^{-1} \quad (6)$$

to the equation (1), we find that (6) really works and relates (1) with the equation

$$v_t = v^3 v_{xxx} + 3v^2 v_x v_{xx} + v^3 v_x. \quad (7)$$

Second, we notice that (7) can be written in the form

$$v_t = v^2 (v v_{xx} + v_x^2 + \frac{1}{2} v^2)_x. \quad (8)$$

Owing to this property, the equation (7) admits the transformation

$$(y, t, w(y, t)) \mapsto (x, t, v(x, t)) : \quad x = w, \quad v = w_y, \quad (9)$$

which turns out to relate (7) with

$$w_t = w_{yyy} + \frac{1}{2} w_y^3. \quad (10)$$

And, third, we make the transformation

$$(y, t, w(y, t)) \mapsto (y, t, z(y, t)) : \quad z = w_y \quad (11)$$

of (10) to the mKdV

$$z_t = z_{yyy} + \frac{3}{2} z^2 z_y, \quad (12)$$

for convenience in what follows, because  $z(y, t) = v(x, t)$ .

### 3 Zero-curvature representation

#### 3.1 Transformation of the mKdV's ZCR

Using the chain of transformations (6), (9) and (11), we can derive a Lax pair for the equation (1), in the form of a ZCR containing an essential parameter, from the following well-known ZCR of the mKdV (12) [4, 5]:

$$\Phi_y = A\Phi, \quad \Phi_t = B\Phi, \quad D_t A = D_y B - [A, B] \quad (13)$$

with

$$A = \begin{pmatrix} \alpha & \frac{i}{2}z \\ \frac{i}{2}z & -\alpha \end{pmatrix}, \quad (14)$$

$$B = \begin{pmatrix} \frac{1}{2}\alpha z^2 + 4\alpha^3 & \frac{i}{2}z_{yy} + \frac{i}{4}z^3 + i\alpha z_y + 2i\alpha^2 z \\ \frac{i}{2}z_{yy} + \frac{i}{4}z^3 - i\alpha z_y + 2i\alpha^2 z & -\frac{1}{2}\alpha z^2 - 4\alpha^3 \end{pmatrix}, \quad (15)$$

where  $\Phi(y, t)$  is a two-component column,  $D_t$  and  $D_y$  stand for the total derivatives,  $[A, B]$  denotes the matrix commutator, and  $\alpha$  is a parameter.

First of all, we obtain a ZCR for the equation (7) through the transformations (9) and (11). Introducing the column  $\Psi : \Psi(x, t) = \Phi(y, t)$ , we have  $\Phi_y = z\Psi_x$ , which allows to rewrite the equation  $\Phi_y = A\Phi$  as

$$\Psi_x = X\Psi, \quad (16)$$

where  $X = z^{-1}A$  after substitution of  $z(y, t) = v(x, t)$ ,

$$X = \begin{pmatrix} \alpha v^{-1} & \frac{i}{2} \\ \frac{i}{2} & -\alpha v^{-1} \end{pmatrix}. \quad (17)$$

The equation  $\Phi_t = B\Phi$ , due to  $\Phi_t = \Psi_t + w_t\Psi_x$  and  $w_t = z_{yy} + \frac{1}{2}z^3$ , leads to

$$\Psi_t = T\Psi, \quad (18)$$

where  $T = B - (z^{-1}z_{yy} + \frac{1}{2}z^2)A$  after substitution of  $z = v$ ,  $z_y = vv_x$  and  $z_{yy} = v^2v_{xx} + vv_x^2$ ,

$$T = \begin{pmatrix} -\alpha(vv_{xx} + v_x^2) + 4\alpha^3 & i\alpha vv_x + 2i\alpha^2 v \\ -i\alpha vv_x + 2i\alpha^2 v & \alpha(vv_{xx} + v_x^2) - 4\alpha^3 \end{pmatrix}. \quad (19)$$

It is easy to check that the compatibility condition

$$D_t X = D_x T - [X, T] \quad (20)$$

of the equations (16) and (18), with the matrices  $X$  (17) and  $T$  (19), determines exactly the equation (7).

Then we can use the transformation (6). Substituting  $v = -(u_{xx} + u)^{-1}$  into  $X$  (17) and  $T$  (19), we obtain a ZCR of the equation (1), in the sense that (20) is satisfied by any solution of (1). This ZCR, however, determines not the equation (1) itself, but a differential prolongation of (1),

$$u_{xxt} + u_t = \frac{1}{2} \left( (u_{xx} + u)^{-2} \right)_{xxx} + \frac{1}{2} \left( (u_{xx} + u)^{-2} \right)_x, \quad (21)$$

due to the structure of the transformed matrix  $X$ ,

$$X = \begin{pmatrix} -\alpha(u_{xx} + u) & \frac{i}{2} \\ \frac{i}{2} & \alpha(u_{xx} + u) \end{pmatrix}. \quad (22)$$

The situation can be improved by a linear transformation of the auxiliary vector function  $\Psi$ ,

$$\Psi \mapsto G\Psi, \quad \det G \neq 0, \quad (23)$$

which generates a gauge transformation of  $X$  and  $T$ ,

$$X \mapsto GXG^{-1} + (D_x G)G^{-1}, \quad T \mapsto GTG^{-1} + (D_t G)G^{-1}. \quad (24)$$

The choice of

$$G = \begin{pmatrix} \exp(\alpha u_x) & 0 \\ 0 & \exp(-\alpha u_x) \end{pmatrix} \quad (25)$$

leads through (24) to the following gauge-transformed matrix  $X$ , which does not contain  $u_{xx}$ :

$$X = \begin{pmatrix} -\alpha u & \frac{i}{2} \exp(2\alpha u_x) \\ \frac{i}{2} \exp(-2\alpha u_x) & \alpha u \end{pmatrix}. \quad (26)$$

Note that  $u$  and  $u_x$  are separated in (26), and a ZCR with such a matrix  $X$  can determine an evolution equation exactly.

Now, from (19), (6), (24) and (25), we obtain the following matrix  $T$ , where  $u_t$  and  $u_{xt}$  have been expressed through (1) in terms of  $x$ -derivatives of  $u$ :

$$T = \begin{pmatrix} 4\alpha^3 & T_{12} \\ T_{21} & -4\alpha^3 \end{pmatrix} \quad (27)$$

with

$$T_{12} = -i\alpha \exp(2\alpha u_x) \left( \frac{2\alpha}{u_{xx} + u} + \frac{u_{xxx} + u_x}{(u_{xx} + u)^3} \right), \quad (28)$$

$$T_{21} = i\alpha \exp(-2\alpha u_x) \left( -\frac{2\alpha}{u_{xx} + u} + \frac{u_{xxx} + u_x}{(u_{xx} + u)^3} \right). \quad (29)$$

It is easy to check that the matrices  $X$  (26) and  $T$  (27)–(29) constitute a ZCR of (1), in the sense that the condition (20) with these matrices determines exactly the equation (1).

### 3.2 Simpler ZCRs are trivial

The obtained ZCR of (1) is characterized by the complicated matrix  $X$  (26) containing  $u_x$ . Does the equation (1) admit any simpler ZCR, with  $X = X(x, t, u)$ , of any dimension  $n \times n$ ? This problem can be solved by direct analysis of the condition (20).

Substituting  $X = X(x, t, u)$  and  $T = T(x, t, u, u_x, u_{xx})$  into (20) and replacing  $u_t$  by the right-hand side of (1), we obtain the following condition, which must be an identity, not an ordinary differential equation restricting solutions of (1):

$$X_t - \frac{u_{xxx} + u_x}{(u_{xx} + u)^3} X_u = D_x T - [X, T] \quad (30)$$

(here and below, subscripts denote derivatives, like  $T_{u_x} = \partial T / \partial u_x$ ). Applying  $\partial / \partial u_{xxx}$  and  $\partial / \partial u_{xx}$  to the identity (30), we obtain, respectively,

$$T_{u_{xx}} = -(u_{xx} + u)^{-3} X_u, \quad (31)$$

$$T_{u_x} = (u_{xx} + u)^{-3} (D_x X_u - [X, X_u]). \quad (32)$$

The compatibility condition  $(T_{u_{xx}})_{u_x} = (T_{u_x})_{u_{xx}}$  for (31) and (32) is  $D_x X_u = [X, X_u]$ , which is equivalent to

$$X = P(x, t)u + Q(x, t) : \quad P_x = [Q, P]. \quad (33)$$

Now, we make use of gauge transformations (24) with  $G = G(x, t)$ , choose  $G$  to be any solution with  $\det G \neq 0$  of the system of ordinary differential equations  $G_x = -GQ$ , and thus set  $Q = 0$  and  $P = P(t)$  in the gauge-transformed matrix  $X$  (33). Then,  $T_u = T_{u_{xx}}$  follows from  $\partial / \partial u_x$  of (30), and this leads through the identity (30) to

$$T = \frac{1}{2} P(t) (u_{xx} + u)^{-2} + K(t) : \quad P_t = [K, P]. \quad (34)$$

Finally, we make  $K = 0$  by a gauge transformation (24) with  $G = G(t)$  satisfying  $G_t = -GK$  and  $\det G \neq 0$ , and thus obtain

$$X = Pu, \quad T = \frac{1}{2}P(u_{xx} + u)^{-2}, \quad P = \text{constant}, \quad (35)$$

with any matrix  $P$  of any dimension  $n \times n$ . However, these matrices  $X$  and  $T$  (35) commute,  $[X, T] = 0$ , and the corresponding ZCR (20) is nothing but  $n^2$  copies of the evident conservation law of the equation (1). In this sense, all the ZCRs sought, with  $X = X(x, t, u)$ , turn out to be trivial, up to gauge transformations (24) with arbitrary  $G(x, t)$ .

## 4 Recursion operator

Let us derive a recursion operator of the equation (1) from the matrix  $X$  (26) of its ZCR. We do it, following the way described in [6] (see also references therein). The recursion operator comes from the problem of finding the class of evolution equations

$$u_t = f(x, t, u, u_x, \dots, u_{x\dots x}) \quad (36)$$

that admit ZCRs (20) with the predetermined matrix  $X$  (26) and any  $2 \times 2$  matrices  $T(\alpha, x, t, u, u_x, \dots, u_{x\dots x})$  of any order in  $u_{x\dots x}$ .

The characteristic form of the ZCR (20) of an equation (36), with  $X$  given by (26), is

$$fC = \nabla_x S, \quad (37)$$

where  $C$  is the characteristic matrix,

$$C = \frac{\partial X}{\partial u} - \nabla_x \left( \frac{\partial X}{\partial u_x} \right), \quad (38)$$

the operator  $\nabla_x$  is defined as  $\nabla_x H = D_x H - [X, H]$  for any (here,  $2 \times 2$ ) matrix  $H$ , and the matrix  $S$  is determined by

$$S = T - f \frac{\partial X}{\partial u_x}. \quad (39)$$

The explicit form of  $C$  (38) for  $X$  (26) is

$$C = \begin{pmatrix} 0 & -2i\alpha^2 e^{2\alpha u_x} (u_{xx} + u) \\ -2i\alpha^2 e^{-2\alpha u_x} (u_{xx} + u) & 0 \end{pmatrix}. \quad (40)$$

Under the gauge transformations (24) with any  $G(\alpha, x, t, u, u_x, \dots, u_{x\dots x})$ , the characteristic matrix  $C$  is transformed as  $C \mapsto GCG^{-1}$  [7], and therefore  $\det C$  is a gauge invariant. We have  $\det C = 4\alpha^4(u_{xx} + u)^2$  in the case of (40), and this proves that the parameter  $\alpha$  cannot be ‘gauged out’ from  $X$  (26), as well as that the matrix  $X$  (26) cannot be transformed by (24) into some  $X$  containing no derivatives of  $u$ .

Computing  $\nabla_x C$ ,  $\nabla_x^2 C$  and  $\nabla_x^3 C$ , we find the cyclic basis to be three-dimensional,  $\{C, \nabla_x C, \nabla_x^2 C\}$ , with the closure equation

$$\nabla_x^3 C = c_0 C + c_1 \nabla_x C + c_2 \nabla_x^2 C, \quad (41)$$

where

$$\begin{aligned} c_0 = & \frac{u_{xxxxx} + 2u_{xxx} + u_x}{u_{xx} + u} \\ & - 9 \frac{(u_{xxx} + u_x)(u_{xxxx} + u_{xx})}{(u_{xx} + u)^2} + 12 \frac{(u_{xxx} + u_x)^3}{(u_{xx} + u)^3}, \end{aligned} \quad (42)$$

$$c_1 = 4 \frac{u_{xxxx} + u_{xx}}{u_{xx} + u} - 12 \frac{(u_{xxx} + u_x)^2}{(u_{xx} + u)^2} + 4\alpha^2(u_{xx} + u)^2 - 1, \quad (43)$$

$$c_2 = 5 \frac{u_{xxx} + u_x}{u_{xx} + u}. \quad (44)$$

Setting  $T$  to be traceless (without loss of generality), we decompose the matrix  $S$  (39) over the cyclic basis as

$$S = s_0 C + s_1 \nabla_x C + s_2 \nabla_x^2 C, \quad (45)$$

where  $s_0$ ,  $s_1$  and  $s_2$  are unknown scalar functions of  $x, t, u, u_x, \dots, u_{x\dots x}$  and  $\alpha$ . Substitution of (45) into (37) leads us through (41) to

$$f = D_x s_0 + c_0 s_2, \quad s_0 = -D_x s_1 - c_1 s_2, \quad s_1 = -D_x s_2 - c_2 s_2, \quad (46)$$

where the function  $s_2$  remains arbitrary. Then, from (46) and (42)–(44), we obtain

$$f = (M - \lambda N) r, \quad (47)$$

where  $\lambda = 4\alpha^2$ ,  $r(\lambda, x, t, u, u_x, \dots, u_{x\dots x}) = s_2$  is any function, of any order

in  $u_{x\dots x}$ , and the linear differential operators  $M$  and  $N$  are

$$\begin{aligned}
M = & D_x^3 + 5 \frac{u_{xxx} + u_x}{u_{xx} + u} D_x^2 \\
& + \left( 6 \frac{u_{xxxx} + u_{xx}}{u_{xx} + u} + 2 \frac{(u_{xxx} + u_x)^2}{(u_{xx} + u)^2} + 1 \right) D_x \\
& + \left( \frac{2u_{xxxxx} + 3u_{xxx} + u_x}{u_{xx} + u} \right. \\
& \left. + 4 \frac{(u_{xxx} + u_x)(u_{xxx} + u_{xx})}{(u_{xx} + u)^2} - 2 \frac{(u_{xxx} + u_x)^3}{(u_{xx} + u)^3} \right), \tag{48}
\end{aligned}$$

$$N = (u_{xx} + u)^2 D_x + 2(u_{xx} + u)(u_{xxx} + u_x). \tag{49}$$

Now, using the expansion

$$r = r_0 + \lambda r_1 + \lambda^2 r_2 + \lambda^3 r_3 + \dots, \tag{50}$$

we obtain from (47) the expression for the right-hand side  $f$  of the represented equation (36), such that  $\partial f / \partial \lambda = 0$  holds,

$$f = Mr_0, \tag{51}$$

as well as the recursion relations for the coefficients  $r_i(x, t, u, u_x, \dots, u_{x\dots x})$  of the expansion (50),

$$Mr_{i+1} = Nr_i, \quad i = 0, 1, 2, \dots \tag{52}$$

The problem has been solved: for any set of functions  $r_0, r_1, r_2, \dots$  satisfying the recursion relations (52), the expression (51) determines an evolution equation (36) admitting a ZCR (20) with the matrix  $X$  given by (26).

It only remains to notice that, if a set of functions  $r_0, r_1, r_2, \dots$  satisfies the recursion relations (52), then the set of functions  $r'_0, r'_1, r'_2, \dots$  determined by  $r'_i = N^{-1}Mr_i$  ( $i = 0, 1, 2, \dots$ ) satisfies (52) as well. Therefore the evolution equation  $u_t = f'$  with  $f' = Mr'_0 = MN^{-1}Mr_0 = MN^{-1}f = Rf$  admits a ZCR (20) with  $X$  (26) whenever an equation  $u_t = f$  does. Eventually, (48) and (49) lead us to the following explicit expression for the recursion operator  $R = MN^{-1}$  of the equation (1):

$$R = \frac{1}{u_{xx} + u} D_x \frac{1}{u_{xx} + u} (D_x + D_x^{-1}). \tag{53}$$



## 5 Conclusion

Some remarks on the obtained results follow.

We succeeded in transforming the new Chou–Qu equation (1) into an integrable equation, the old and well-studied mKdV. The applicability of Miura-type transformations, however, is not restricted to integrable equations only. For instance, the original Miura transformation relates very wide (continual) classes of (mainly non-integrable) evolution equations [8].

We found the simplest nontrivial ZCR of the evolution equation (1). Its matrix  $X$  (26) contains  $u_x$ . For this reason, such a ZCR cannot be detected by those existent methods, which assume, as a starting point, that  $X = X(x, t, u)$  must suffice in the case of evolution equations.

Of course, we could derive the obtained recursion operator (53) from the well-known recursion operator of the mKdV through the transformations found. However, we used a different method instead, mainly in order to illustrate, by this rather complicated example of  $X = X(\alpha, u, u_x)$  (26), how the method works algorithmically.

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